

Gibbs Equation

1 Overview

Recall the second law for CM, we can relate entropy changes to heat transfer and irreversibility:

$$dS = \frac{\delta Q}{T} + \delta P_s \quad (1)$$

But entropy is a TD property, we want to relate S changes to changes in other TD properties. Therefore, we need to find a generally applicable state equation.

2 Gibbs Equation

2.1 Derivation

Recall the first law:

$$de = \delta q + \delta w \quad (2)$$

If we only consider internal energy, and the only reversible work mode is compression/expansion:

$$du = \delta q - p dv \quad (3)$$

If we assume the process is reversible, based on second law:

$$ds = \delta q/T \quad (4)$$

Plug back in, we get:

$$du = T ds - p dv \quad (5)$$

Solve ds , we get the **Gibbs Equation**:

$$ds = \frac{du}{T} + \frac{p}{T} dv \quad (6)$$

Recall the definition of enthalpy:

$$dh = du + p dv + v dp \quad (7)$$

We get another version of **Gibbs Equation**:

$$ds = \frac{dh}{T} - \frac{v}{T}dp \quad (8)$$

2.2 State Equations

During the derivation, we assume **reversible process with only compression/expansion work**. However, these expressions **are valid not only under these conditions, but generally applicable**, since they only involve **state properties**, they are **state equations**.

2.3 For Ideal Gases

Start from original Gibbs Equation:

$$ds = \frac{du}{T} + \frac{p}{T}dv \quad (9)$$

Recall that:

$$du = c_v dT \quad (10)$$

$$pv = RT \quad (11)$$

So we have:

$$ds = \frac{c_v dT}{T} + \frac{R}{v}dv \quad (12)$$

If we integrate this:

$$\int_{s_1}^{s_2} ds = \int_{T_1}^{T_2} \frac{c_v dT}{T} + \int_{v_1}^{v_2} \frac{R}{v} dv \quad (13)$$

$$s_2 - s_1 = \int_{T_1}^{T_2} \frac{c_v dT}{T} + R \int_{v_1}^{v_2} \frac{dv}{v} \quad (14)$$

Finally we have:

$$s_2 - s_1 = \int_{T_1}^{T_2} \frac{c_v dT}{T} + R \ln \frac{v_2}{v_1} \quad (15)$$

Or:

$$s_2 - s_1 = \int_{T_1}^{T_2} \frac{c_v dT}{T} - R \ln \frac{\rho_2}{\rho_1} \quad (16)$$

Similarly, we can use enthalpy version:

$$ds = \frac{dh}{T} - \frac{v}{T}dp = \frac{c_p dT}{T} - \frac{R}{p}dp \quad (17)$$

Integrate:

$$\int_{s_1}^{s_2} ds = \int_{T_1}^{T_2} \frac{c_p dT}{T} - R \int_{p_1}^{p_2} \frac{dp}{p} \quad (18)$$

$$s_2 - s_1 = \int_{T_1}^{T_2} \frac{c_p dT}{T} - R \ln \frac{p_2}{p_1} \quad (19)$$

Notice that the dependencies of entropy respect to T , p and ρ are separate.

2.4 For Carlantically Perfect Gas

Under CPG assumption, specific heat is constant. Therefore, we have:

$$s_2 - s_1 = c_v \ln \frac{T_2}{T_1} + R \ln \frac{v_2}{v_1} \quad (20)$$

$$s_2 - s_1 = c_v \ln \frac{T_2}{T_1} - R \ln \frac{\rho_2}{\rho_1} \quad (21)$$

$$s_2 - s_1 = c_p \ln \frac{T_2}{T_1} - R \ln \frac{p_2}{p_1} \quad (22)$$

Normalization:

$$\frac{\Delta s_{12}}{R} = \frac{c_p}{R} \ln \frac{T_2}{T_1} - \ln \frac{p_2}{p_1} = \ln \left(\frac{T_2}{T_1} \right)^{c_p/R} - \ln \frac{p_2}{p_1} \quad (23)$$

Recall that:

$$\frac{c_p}{R} = \frac{\gamma}{\gamma - 1} \quad (24)$$

$$\frac{c_v}{R} = \frac{1}{\gamma - 1} \quad (25)$$

Therefore:

$$\frac{\Delta s_{12}}{R} = \ln \left(\frac{T_2}{T_1} \right)^{\frac{\gamma}{\gamma-1}} - \ln \frac{p_2}{p_1} = \ln \left(\frac{T_2}{T_1} \right)^{\frac{1}{\gamma-1}} - \ln \frac{p_2}{p_1} \quad (26)$$

3 Isentropic Relations

3.1 TPG

At first, we assume **TPG**, which means specific heat is a function of temperature. For simplification, we define:

$$s_1^o = \int_{T_{ref}}^{T_1} \frac{c_p dT}{T} \quad (27)$$

$$s_2^o = \int_{T_{ref}}^{T_2} \frac{c_p dT}{T} \quad (28)$$

$$\Delta s_{12}^o(T) = s_2^o - s_1^o \quad (29)$$

For isentropic process, $\Delta s = 0$. Therefore:

$$0 = \int_{T_1}^{T_2} \frac{c_p dT}{T} - R \ln \frac{p_2}{p_1} \quad (30)$$

$$\ln \frac{p_2}{p_1} = \frac{1}{R} \int_{T_1}^{T_2} \frac{c_p dT}{T} \quad (31)$$

Finally we have:

$$\frac{p_2}{p_1} = e^{\frac{1}{R} \int_{T_1}^{T_2} \frac{c_p dT}{T}} = e^{\frac{\Delta s_{12}^o(T)}{R}} \quad (32)$$

In terms of specific volume, recall that:

$$\frac{v_1}{v_2} = \frac{T_1/p_1}{T_2/p_2} = \frac{T_1 p_2}{T_2 p_1} \quad (33)$$

Therefore:

$$\frac{v_1}{v_2} = \frac{\rho_2}{\rho_1} = \frac{T_1}{T_2} e^{\frac{\Delta s_{12}^o(T)}{R}} \quad (34)$$

3.2 CPG

Now, the specific heat is a constant, then:

$$0 = c_p \ln \frac{T_2}{T_1} - R \ln \frac{p_2}{p_1} \quad (35)$$

$$\frac{p_2}{p_1} = \left(\frac{T_2}{T_1} \right)^{\frac{c_p}{R}} = \left(\frac{T_2}{T_1} \right)^{\frac{\gamma}{\gamma-1}} \quad (36)$$

$$0 = c_v \ln \frac{T_2}{T_1} - R \ln \frac{\rho_2}{\rho_1} \quad (37)$$

$$\frac{\rho_2}{\rho_1} = \left(\frac{T_2}{T_1}\right)^{\frac{c_v}{R}} = \left(\frac{T_2}{T_1}\right)^{\frac{1}{\gamma-1}} \quad (38)$$

Therefore we have the power law of isentropic process:

$$\frac{\rho_2}{\rho_1} = \left(\frac{p_2}{p_1}\right)^{\frac{1}{\gamma}} \quad (39)$$

$$\frac{p}{\rho^\gamma} = pv^\gamma = \text{const} \quad (40)$$

3.3 Polytropic Process

We define the general form as polytropic process:

$$pv^n = \text{constant} \quad (41)$$

Then we have the following conditions:

1. $n = 0$: Isobaric
2. $n = 1$: Isothermal
3. $n = \gamma$: Isentropic
4. $n = \infty$: Isochoric